

# 球面上のベクトル場の推定とその応用

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## Localized basis functions for spherical vector field

- A divergence-free vector field can be represented by a stream function  $\Psi^{df}$  as follows:

$$\mathbf{V}^{df}(\mathbf{r}) = -\mathbf{e}_r \times \nabla \Psi^{df}.$$

- A curl-free vector field can be represented by a potential function  $\Psi^{cf}$  as follows:

$$\mathbf{V}^{cf}(\mathbf{r}) = -\nabla \Psi^{cf}.$$

- According to the Helmholtz theorem, an arbitrary vector field on a sphere can be written as a sum of a divergence-free field and a curl-free field. We can thus represent any vector field in the following form:

$$\mathbf{V}(\mathbf{r}) = \mathbf{V}^{df}(\mathbf{r}) + \mathbf{V}^{cf}(\mathbf{r}) = -\mathbf{e}_r \times \nabla \Psi^{df} - \nabla \Psi^{cf}.$$

- We expand the stream function  $\Psi^{df}$  and the potential function  $\Psi^{cf}$  by using localized basis functions  $\psi^{df}$  and  $\psi^{cf}$ :

$$\Psi^{df}(\mathbf{r}) = \sum_i w_i^{df} \psi^{df}(\mathbf{r}, \mathbf{r}_i), \quad \Psi^{cf}(\mathbf{r}) = \sum_i w_i^{cf} \psi^{cf}(\mathbf{r}, \mathbf{r}_i).$$

- Defining the following vector-valued basis functions:

$$\mathbf{v}^{df}(\mathbf{r}, \mathbf{r}_i) = -\mathbf{e}_r \times \nabla \psi^{df}(\mathbf{r}, \mathbf{r}_i), \quad \mathbf{v}^{cf}(\mathbf{r}, \mathbf{r}_i) = -\nabla \psi^{cf}(\mathbf{r}, \mathbf{r}_i),$$

an arbitrary spherical field can be represented using these basis functions

$$\begin{aligned} \mathbf{V}(\mathbf{r}) &= \sum_i w_i^{df} \left( -\mathbf{e}_r \times \nabla \psi^{df}(\mathbf{r}, \mathbf{r}_i) \right) + \sum_i w_i^{cf} \left( -\nabla \psi^{cf}(\mathbf{r}, \mathbf{r}_i) \right) \\ &= \sum_i w_i^{df} \mathbf{v}^{df}(\mathbf{r}, \mathbf{r}_i) + \sum_i w_i^{cf} \mathbf{v}^{cf}(\mathbf{r}, \mathbf{r}_i). \end{aligned}$$

## Modeling of divergence-free vector field

- We employ spherical Gaussian functions for obtaining a set of basis functions

$$\psi^{df}(\mathbf{r}, \mathbf{r}_i) = \psi^{cf}(\mathbf{r}, \mathbf{r}_i) = \exp \left[ \eta \left( \frac{\mathbf{r} \cdot \mathbf{r}_i}{R^2} - 1 \right) \right] = \exp \left[ \eta (\cos \theta' - 1) \right].$$

- We then obtain

$$\mathbf{v}^{df}(\mathbf{r}, \mathbf{r}_i) = (\eta \mathbf{r}_i \times \mathbf{r}) \exp \left[ \eta \left( \frac{\mathbf{r} \cdot \mathbf{r}_i}{R^2} - 1 \right) \right], \quad \mathbf{v}^{cf}(\mathbf{r}, \mathbf{r}_i) = \mathbf{e}_{\theta, i} (\eta \sin \Delta \theta) \exp \left[ \eta \left( \frac{\mathbf{r} \cdot \mathbf{r}_i}{R^2} - 1 \right) \right].$$

- Now we consider a divergence-free vector field (no source, no sink) and expand the field by using the divergence-free basis functions:

$$\mathbf{V}(\mathbf{r}) = \sum_i w_i^{df} \mathbf{v}^{df}(\mathbf{r}, \mathbf{r}_i).$$

- When we use the expansion of our basis functions

$$\mathbf{V}(\mathbf{r}) = \sum_i w_i \mathbf{v}(\mathbf{r}, \mathbf{r}_i),$$

the node points  $\mathbf{r}_i$  can be placed arbitrarily.

- We placed 2500 node points randomly and uniformly distributed in the region above 40 degree in latitude, which approximates the following Monte Carlo convolution:

$$\mathbf{V}(\mathbf{r}) = \int_S w(\mathbf{r}') \mathbf{v}(\mathbf{r}, \mathbf{r}') d\mathbf{r}' = \int_0^{\pi/2} \int_0^{2\pi} w(\mathbf{r}') \mathbf{v}(\mathbf{r}, \mathbf{r}') R^2 \sin \theta d\theta d\phi.$$

## Application to ionospheric physics

- Now we apply this method for estimating the ionospheric plasma velocity distribution which can be assumed to be divergence-free.
- We fit the model to the data of SuperDARN, which is a radar network observing the ionospheric plasma velocity.
- The gaps of the data coverage of SuperDARN are filled with the empirical model (Weimber 2001).
- We assume the weight  $w$  can be decomposed into the model-based value  $\zeta$  and the residual  $\beta$ :

$$w = \zeta + \beta,$$

and the residual  $\beta$  is estimated with the Kalman filter.

- We assume the temporal evolution of the weights  $\beta$

$$p(\beta_k | \beta_{k-1}) = \mathcal{N}(\alpha \beta_{k-1}, Q).$$

- The residual component can then be estimated with the following Kalman filter algorithm:

Prediction:

$$\beta_{k|k-1} = \alpha \beta_{k-1|k-1},$$

$$P_{k|k-1} = \alpha^2 P_{k-1|k-1} + Q.$$

Filtering:

$$\beta_{k|k} = \beta_{k|k-1} + (P_{k|k-1}^{-1} + H_k^T R_k^{-1} H_k)^{-1} H_k^T R_k^{-1} (y_k - H_k \beta_{k|k-1}),$$

$$P_{k|k} = (P_{k|k-1}^{-1} + H_k^T R_k^{-1} H_k)^{-1}.$$

## Design of covariance matrices

- In order to ensure spatial smoothness, the covariance matrix  $Q_k$  is set as follows:

$$Q = \sigma_Q^2 \begin{pmatrix} C_Q(\mathbf{r}_1, \mathbf{r}_1) & \cdots & C_Q(\mathbf{r}_1, \mathbf{r}_n) \\ \vdots & \ddots & \vdots \\ C_Q(\mathbf{r}_n, \mathbf{r}_1) & \cdots & C_Q(\mathbf{r}_n, \mathbf{r}_n) \end{pmatrix},$$

where  $C_Q(\mathbf{r}_i, \mathbf{r}_j)$  is the following correlation function:

$$C_Q(\mathbf{r}_i, \mathbf{r}_j) = \rho(\mathbf{r}_i, \mathbf{r}_j) \exp \left[ \kappa \left( \frac{\mathbf{r}_i \cdot \mathbf{r}_j}{R^2} - 1 \right) \right],$$

$$\text{where } \rho(\mathbf{r}_i, \mathbf{r}_j) = \frac{(\sin \lambda_i - \sin 40^\circ)(\sin \lambda_j - \sin 40^\circ)}{(1 - \sin 40^\circ)^2}.$$

- The matrix  $R_k$  is set as follows:

$$R_k = \sigma_R^2 \begin{pmatrix} \delta_{b_1 b_1} C_R(g_1, g_1) & \cdots & \delta_{b_1 b_1} C_R(g_1, g_1) \\ \vdots & \ddots & \vdots \\ \delta_{b_1 b_1} C_R(g_1, g_1) & \cdots & \delta_{b_1 b_1} C_R(g_1, g_1) \end{pmatrix},$$

where  $b_i$  and  $g_i$  denote the beam number and the range gate of the  $i$ -th element of the observation  $y$  and

$$C_R(g_i, g_j) = \exp \left[ -\frac{(g_i - g_j)^2}{2} \right].$$

Results >> See our paper for detail ( <https://doi.org/10.1186/s40623-020-01168-4> ).